# Bethe-Salpeter equation for the particle-particle propagator

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December 13, 2024



This project has received funding from the European Research Council (ERC) under the European Union's Horizon 2020 research and innovation programme (Grant agreement No. 863481).

# Definition $\begin{array}{c} \underbrace{1 = (\mathbf{r}_1, \sigma_1, t_1)}_{G_2(12; 1'2') = (-\mathbf{i})^2 \left\langle \begin{array}{c} \Psi_0^{\mathsf{N}} \\ \end{array} \middle| \hat{T} [ \hat{\psi}(1) \quad \hat{\psi}(2) \quad \hat{\psi}^{\dagger}(2') \quad \hat{\psi}^{\dagger}(1') \end{array} \right| \left| \begin{array}{c} \Psi_0^{\mathsf{N}} \\ \end{array} \right\rangle}_{\overrightarrow{\mathsf{N}-\text{electron ground-state}} \\ \overrightarrow{\mathsf{Field operators}} \end{array}$

## • Electron-hole pair propagation

$$\mathbf{t}_{2}, \mathbf{t}_{2'} > \mathbf{t}_{1}, \mathbf{t}_{1'} \qquad G_{2}(12; 1'2') = (-\mathrm{i})^{2} \left\langle \Psi_{0}^{N} \right| \hat{\mathcal{T}} \Big[ \hat{\psi}(2) \hat{\psi}^{\dagger}(2') \Big] \hat{\mathcal{T}} \Big[ \hat{\psi}(1) \hat{\psi}^{\dagger}(1') \Big] \big| \Psi_{0}^{N} \right\rangle$$

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## • Electron-electron pair propagation

$$\mathbf{t}_{1}, \mathbf{t}_{2} > \mathbf{t}_{1'}, \mathbf{t}_{2'} \qquad G_{2}(12; 1'2') = (-\mathrm{i})^{2} \left\langle \Psi_{0}^{N} \middle| \hat{\mathcal{T}} \Big[ \hat{\psi}(1) \hat{\psi}(2) \Big] \hat{\mathcal{T}} \Big[ \hat{\psi}^{\dagger}(1') \hat{\psi}^{\dagger}(2') \Big] \middle| \Psi_{0}^{N} \right\rangle$$

• Hole-hole pair propagation

$$\mathbf{t}_{1'}, \mathbf{t}_{2'} > \mathbf{t}_1, \mathbf{t}_2 \qquad G_2(12; 1'2') = (-\mathrm{i})^2 \left\langle \Psi_0^N \right| \hat{\mathcal{T}} \Big[ \hat{\psi}^{\dagger}(1') \hat{\psi}^{\dagger}(2') \Big] \hat{\mathcal{T}} \Big[ \hat{\psi}(1) \hat{\psi}(2) \Big] \big| \Psi_0^N \right\rangle$$

#### **Electron-hole propagator**

$$L(12; 1'2') = -G_2(12; 1'2') + G(11')G(22')$$

#### Lehman representation

$$\underbrace{X_{1'} = (\mathbf{r}_{1'}, \sigma_{1'})}_{L(\mathbf{x}_{1}\mathbf{x}_{2}; \mathbf{x}_{1'}\mathbf{x}_{2'}; \omega)} = \sum_{\nu > 0} \frac{L_{\nu}^{N}(\mathbf{x}_{2}\mathbf{x}_{2'})R_{\nu}^{N}(\mathbf{x}_{1}\mathbf{x}_{1'})}{\omega - (E_{\nu}^{N} - E_{0}^{N}) + i\eta} - \sum_{\nu > 0} \frac{L_{\nu}^{N}(\mathbf{x}_{2}\mathbf{x}_{2'})R_{\nu}^{N}(\mathbf{x}_{1}\mathbf{x}_{1'})}{\omega - (E_{0}^{N} - E_{\nu}^{N}) - i\eta}$$

$$\underbrace{N-\text{th Excitation energies}}$$

#### **Electron-hole Bethe-Salpeter equation**

$$L(12; 1'2') = L_0(12; 1'2') + \int d(3456) L_0(14; 1'3) \Xi^{eh}(36; 45) L(52; 62')$$
Independent particle propagator
eh kernel

where 
$$L_0(12; 1'2') = G(12')G(21')$$
  $\Xi^{eh}(12; 34) = \left. \frac{\delta \Sigma(13)}{\delta G(42)} \right|_{U=0}$ 



#### **Eigenvalue problem**

$$egin{pmatrix} oldsymbol{A} & oldsymbol{B} \ oldsymbol{B}^\dagger & oldsymbol{A} \end{pmatrix}egin{pmatrix} oldsymbol{X} \ oldsymbol{Y} \end{pmatrix} = \omega egin{pmatrix} oldsymbol{1} & oldsymbol{0} \ oldsymbol{0} & -oldsymbol{1} \end{pmatrix}egin{pmatrix} oldsymbol{X} \ oldsymbol{Y} \end{pmatrix}$$

$$A_{ia,jb} = (\epsilon_a - \epsilon_i)\delta_{ab}\delta_{ij} + \Xi^{eh}_{ia,jb}$$
$$B_{ia,bj} = \Xi^{eh}_{ia,bj}$$

#### **Eigenvalue problem**

$$egin{pmatrix} oldsymbol{A} & oldsymbol{B} \ oldsymbol{B}^\dagger & oldsymbol{A} \end{pmatrix}egin{pmatrix} oldsymbol{X} \ oldsymbol{Y} \end{pmatrix} = \omegaegin{pmatrix} oldsymbol{1} & oldsymbol{0} \ oldsymbol{0} & -oldsymbol{1} \end{pmatrix}egin{pmatrix} oldsymbol{X} \ oldsymbol{Y} \end{pmatrix}$$

#### Kernels

$$\Xi_{ia,jb}^{\text{eh,RPA}} = \langle ib|aj \rangle \qquad \Xi_{ia,jb}^{\text{eh,GW}} = \langle ib|aj \rangle - W_{ibja}(\omega = 0)$$

$$\begin{aligned} \mathsf{A}_{ia,jb} &= (\epsilon_a - \epsilon_i) \delta_{ab} \delta_{ij} + \Xi^{\mathsf{eh}}_{ia,jb} \\ \mathsf{B}_{ia,bj} &= \Xi^{\mathsf{eh}}_{ia,bj} \end{aligned}$$

- Second-order kernel
- *T*-matrix kernel

• ...

#### **Schwinger relation**

$$-G_{2}(12; 1'2') + G(11')G(22') = \frac{\delta G(11'; [U])}{\delta U^{eh}(2'2)} \bigg|_{U=0} = L(12; 1'2')$$
  
External potential

**External potential** 

$$\hat{\mathcal{U}}(t_1) = \int \mathrm{d}(\mathbf{x}_1 \mathbf{x}_{1'} t_1') \, \hat{\psi}^{\dagger}(\mathbf{x}_1) \boldsymbol{U}^{\mathsf{eh}}(\mathbf{11'}) \hat{\psi}(\mathbf{x}_{1'})$$

#### Another external potential ...

$$\hat{\mathcal{U}}(t_1) = \frac{1}{2} \left( \int \mathrm{d}(\mathbf{x}_1 \mathbf{x}_{1'} t_1') \, \hat{\psi}(\mathbf{x}_1) \mathcal{U}^{\mathsf{h}\mathsf{h}}(11') \hat{\psi}(\mathbf{x}_{1'}) + \int \mathrm{d}(\mathbf{x}_1 \mathrm{d}\mathbf{x}_{1'} t_1') \, \hat{\psi}^{\dagger}(\mathbf{x}_{1'}) \, \mathcal{U}^{\mathsf{ee}}(11') \, \hat{\psi}^{\dagger}(\mathbf{x}_{1'}) \right)$$
Non-number conserving

#### Another external potential ...

$$\hat{\mathcal{U}}(t_1) = \frac{1}{2} \left( \int \mathrm{d}(\mathbf{x}_1 \mathbf{x}_{1'} t_1') \, \hat{\psi}(\mathbf{x}_1) U^{\mathsf{h}\mathsf{h}}(11') \hat{\psi}(\mathbf{x}_{1'}) + \int \mathrm{d}(\mathbf{x}_1 \mathrm{d}\mathbf{x}_{1'} t_1') \, \hat{\psi}^{\dagger}(\mathbf{x}_{1'}) \, U^{\mathsf{ee}}(11') \, \hat{\psi}^{\dagger}(\mathbf{x}_{1'}) \right) \\ \mathsf{Non-number conserving}$$

#### ...leading to an alternative Schwinger relation

$$\frac{1}{2} \left( -G_2(12; 1'2'; [U]) + G^{\mathsf{hh}}(12; [U]) G^{\mathsf{ee}}(2'1'; [U]) \right) \bigg|_{U=0} = \left. \frac{\delta G^{\mathsf{ee}}(1'2'; [U])}{\delta U^{\mathsf{hh}}(12)} \right|_{U=0} = \mathcal{K}(12; 1'2')$$

#### **Anomalous propagators**

$$G^{\mathsf{h}\mathsf{h}}(11';[U]) = (-\mathrm{i}) \langle \Psi_0 | \hat{\mathcal{T}} \Big[ \hat{\psi}(1) \hat{\psi}(1') \Big] | \Psi_0 \rangle \quad G^{\mathsf{e}\mathsf{e}}(11';[U]) = (-\mathrm{i}) \langle \Psi_0 | \hat{\mathcal{T}} \Big[ \hat{\psi}^\dagger(1) \hat{\psi}^\dagger(1') \Big] | \Psi_0 \rangle$$

#### Nambu formalism and the Gorkov propagator

$$\mathbf{G}(11') = (-i) \langle \Psi_0 | \hat{\mathcal{T}} \left[ \begin{pmatrix} \hat{\psi}(1) \hat{\psi}^{\dagger}(1') & \hat{\psi}(1) \hat{\psi}(1') \\ \hat{\psi}^{\dagger}(1) \hat{\psi}^{\dagger}(1') & \hat{\psi}^{\dagger}(1) \hat{\psi}(1') \end{pmatrix} \right] | \Psi_0 \rangle = \begin{pmatrix} G^{\mathsf{he}}(11') & G^{\mathsf{hh}}(11') \\ G^{\mathsf{ee}}(11') & G^{\mathsf{eh}}(11') \end{pmatrix}.$$

**Gorkov-Dyson equation** 

$$\mathbf{G}^{-1}(11') = \mathbf{G}_0^{-1}(11') - \begin{pmatrix} \Sigma^{\text{he}}(11') & \Sigma^{\text{hh}}(11') + U^{\text{ee}}(11') \\ \Sigma^{\text{ee}}(11') + U^{\text{hh}}(11') & \Sigma^{\text{eh}}(11') \end{pmatrix}$$

#### Particle-particle Bethe-Salpeter equation

$$K(12;1'2') = K_0(12;1'2') - \int d(3456) K(12;56) \Xi^{pp}(56;34) K_0(34;1'2')$$
Independent particle propagator
pp kernel

where 
$$K_0(12; 1'2') = \frac{1}{2} [G(21')G(12') - G(11')G(22')] = \Xi^{pp}(56; 34) = \left. \frac{\delta \Sigma^{ee}(34)}{\delta G^{ee}(56)} \right|_{U=0}$$



#### Lehman representation

$$\mathcal{K}(\mathbf{x}_{1}\mathbf{x}_{2};\mathbf{x}_{1'}\mathbf{x}_{2'};\omega) = \sum_{\nu} \frac{L_{\nu}^{N+2}(\mathbf{x}_{1}\mathbf{x}_{2})R_{\nu}^{N+2}(\mathbf{x}_{1}'\mathbf{x}_{2}')}{\omega - (E_{\nu}^{N+2} - E_{0}^{N}) + i\eta} - \sum_{\nu} \frac{L_{\nu}^{N-2}(\mathbf{x}_{1}'\mathbf{x}_{2}')R_{\nu}^{N-2}(\mathbf{x}_{1}\mathbf{x}_{2})}{\omega - (E_{0}^{N} - E_{\nu}^{N-2}) - i\eta}$$
Double electron affinities
Double ionization potentials

#### **Eigenvalue problem**

$$egin{pmatrix} m{C} & m{B} \ m{B}^\dagger & m{D} \end{pmatrix} egin{pmatrix} m{X} \ m{Y} \end{pmatrix} = \omega egin{pmatrix} 1 & 0 \ 0 & -1 \end{pmatrix} egin{pmatrix} m{X} \ m{Y} \end{pmatrix}$$

$$\begin{aligned} C^{\text{RPA}}_{ab,cd} &= (\epsilon_a + \epsilon_b) \delta_{ac} \delta_{bd} + \Xi^{\text{PP}}_{ab,cd} \\ B^{\text{RPA}}_{ab,ij} &= \Xi^{\text{PP}}_{ab,ij} \\ D^{\text{RPA}}_{ij,kl} &= -(\epsilon_i + \epsilon_j) \delta_{ik} \delta_{jl} + \Xi^{\text{PP}}_{ij,kl} \end{aligned}$$

#### Eigenvalue problem

$$egin{pmatrix} m{C} & m{B} \ m{B}^\dagger & m{D} \end{pmatrix} egin{pmatrix} m{X} \ m{Y} \end{pmatrix} = \omega egin{pmatrix} 1 & 0 \ 0 & -1 \end{pmatrix} egin{pmatrix} m{X} \ m{Y} \end{pmatrix}$$

#### Kernels

$$\Xi_{ij,kl}^{\text{pp,RPA}} = \langle ij||kl \rangle \quad \Xi_{ij,kl}^{\text{pp,GW}} = W_{ijkl}(\omega = 0) - W_{ijlk}(\omega = 0)$$

$$\begin{aligned} C^{\text{RPA}}_{ab,cd} &= (\epsilon_a + \epsilon_b) \delta_{ac} \delta_{bd} + \Xi^{\text{PP}}_{ab,cd} \\ B^{\text{RPA}}_{ab,ij} &= \Xi^{\text{PP}}_{ab,ij} \\ D^{\text{RPA}}_{ij,kl} &= -(\epsilon_i + \epsilon_j) \delta_{ik} \delta_{jl} + \Xi^{\text{PP}}_{ij,kl} \end{aligned}$$

- Second-order kernel
- *T*-matrix kernel

• ...







Sangalli et al. J. Chem. Phys. 158 (2011) 034115

#### Error with respect to CVS-FCI in the aug-cc-pCVTZ basis set



## Conclusion and open questions

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- Simple expression for the kernel of the particle-particle channel
- ppBSE brings quantitative improvements for double ionization
- More details in arXiv:2411.13167

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## **Open questions/problems**

- Correlation energy through adiabatic connection
- Auger spectrum at the GW level



Extracted from Jayadev et al. J. Chem. Phys. 158 (2023) 064109

## **Questions?**

#### Derivation

$$\begin{split} \mathcal{K}(12;1'2') &= \left. \frac{\delta G^{\text{ee}}(1'2')}{\delta U^{\text{hh}}(12)} \right|_{U=0} \tag{1} \\ &= G(31') \left. \frac{\delta (G^{-1})^{\text{ee}}(33')}{\delta U^{\text{hh}}(12)} \right|_{U=0} G(3'2') \tag{2} \\ &= -G(31') \left. \frac{\delta U^{\text{hh}}(33')}{\delta U^{\text{hh}}(12)} \right|_{U=0} G(3'2') - G(31') \left. \frac{\delta \Sigma^{\text{ee}}(33')}{\delta U^{\text{hh}}(12)} \right|_{U=0} G(3'2') \tag{3} \\ &= \mathcal{K}_0(12;1'2') - \left. \frac{\delta G^{\text{ee}}(44')}{\delta U^{\text{hh}}(12)} \right|_{U=0} \left. \frac{\delta \Sigma^{\text{ee}}(33')}{\delta G^{\text{ee}}(44')} \right|_{U=0} G(3'2')G(31') \tag{4} \end{split}$$

#### Self-energy

$$\boldsymbol{\Sigma}(11') = i \int d(33') \begin{pmatrix} W(13'; 31') G^{he}(33') & -W(13'; 31') G^{hh}(33') \\ -W(31'; 13') G^{ee}(33') & W(31'; 13') G^{eh}(33') \end{pmatrix}$$

#### **Screened interaction**



#### eh GW kernel

$$\begin{split} \Xi^{\mathsf{eh},GW}(12;1'2') &= \mathrm{i} \left. \frac{\delta \Sigma(11')}{\delta G(2'2)} \right|_{U=0} = - \left. \frac{\delta \left[ G(33')W(11';33') \right]}{\delta G(2'2)} \right|_{U=0} \\ &= -W(11';2'2) - G(33') \left. \frac{\delta W(11';33')}{\delta G(2'2)} \right|_{U=0} \end{split}$$

#### pp GW kernel

$$\mathbf{i}\Xi^{\mathsf{ee},\mathsf{GW}}(11';22') = \mathbf{i} \left. \frac{\delta\Sigma^{\mathsf{ee}}(22')}{\delta \mathsf{G}^{\mathsf{ee}}(11')} \right|_{U=0} = \left. \frac{\delta \left[ \mathsf{G}^{\mathsf{ee}}(33') \mathcal{W}(33';22') \right]}{\delta \mathsf{G}^{\mathsf{ee}}(11')} \right|_{U=0} = \frac{1}{2} \left[ \mathcal{W}(11';22') - \mathcal{W}(11;2'2) \right] = \frac{1}{2} \left[ \mathcal{W}(11';22') - \mathcal{W}(11;2'2) \right]$$



