Anomalous propagators and the particle-particle correlation channel of many-body perturbation theory: Hedin's equations

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Definitions, Hedin's equations and usual approximations

One-body Green's function

Definition

$1 = (\mathbf{r}_1, \sigma_1, t_1)$

$$G(11') = (-\mathrm{i}) \left\langle \Psi_0^N \middle| \hat{T}\left[\hat{\psi}(1) \; \hat{\psi}^\dagger(1') \right] \middle| \Psi_0^N \right\rangle$$

Field operators

N-electron ground-state

One-body Green's function

Definition

$$G(11') = (-\mathrm{i}) \left\langle \Psi_0^N \middle| \hat{T} \left[\hat{\psi}(1) \hat{\psi}^\dagger(1') \right] \middle| \Psi_0^N \right\rangle$$

Schrödinger picture

$$G(11') = (-\mathrm{i}) \left(\underbrace{\left\langle \Psi_0^N(t_1) \middle| \hat{\psi}(\mathbf{r}_1, \sigma_1) \hat{U}(t_1, t_{1'}) \hat{\psi}^\dagger(\mathbf{r}_{1'}, \sigma_{1'}) \middle| \Psi_0^N(t_{1'}) \right\rangle}_{\text{Propag. of an electron}} \right.$$
$$\left. - \underbrace{\left\langle \Psi_0^N(t_{1'}) \middle| \hat{\psi}^\dagger(\mathbf{r}_{1'}, \sigma_{1'}) \hat{U}(t_{1'}, t_1) \hat{\psi}(\mathbf{r}_1, \sigma_1) \middle| \Psi_0^N(t_1) \right\rangle}_{\text{Propag. of a hole}} \right)$$

Charged excitations

Lehmann representation

$x_1 = (\mathbf{r}_1, \sigma_1)$

$$G(\mathbf{x}_1 \mathbf{x}_{1'} ; \omega) = \sum_S \frac{\mathcal{I}_S(\mathbf{x}_1)\mathcal{I}_S^*(\mathbf{x}_{1'})}{\omega - (E_0^N - E_S^{N-1}) - i\eta} + \sum_S \frac{\mathcal{A}_S(\mathbf{x}_1)\mathcal{A}_S^*(\mathbf{x}_{1'})}{\omega - (E_S^{N+1} - E_0^N) + i\eta}$$

S-th ionization potentials

S-th electron affinities

Reduced quantity theories

Link to one-body density

$$\rho(\mathbf{r}_1) = \lim_{t_{1'} \to t_1} \lim_{\mathbf{r}_{1'} \to \mathbf{r}_1} \sum_{\sigma_1, \sigma_{1'}} G(\mathbf{x}_1, \mathbf{x}_{1'};\ t_{1'} - t_1\ )$$

$\hat{H}$ time-indep.

Reduced quantity theories

Link to one-body density

$$\rho(\mathbf{r}_1) = \lim_{t_{1'} \to t_1} \lim_{\mathbf{r}_{1'} \to \mathbf{r}_1} \sum_{\sigma_1, \sigma_{1'}} G(\mathbf{x}_1, \mathbf{x}_{1'};\ t_{1'} - t_1\ )$$

$\hat{H}$ time-indep.

Spectral function

$$A(\omega) = \frac{1}{\pi} |\text{Im}\, G(\omega)|$$

Total energy

$$E^{\text{GM}} = -\frac{\mathrm{i}}{2} \int d\mathbf{x}_1 \lim_{2 \to 1^+} \left[\mathrm{i} \frac{\partial}{\partial t_1} + h(\mathbf{x}_1) \right] G(12)$$

How to compute G?

The Dyson equation

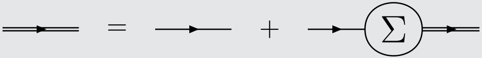
$$G(11') = G_0(11') + \int \mathrm{d}(22')\, G_0(12)\, \Sigma(22')\, G(2'1')$$

Self-energy

How to compute G?

The Dyson equation

$$G(11') = G_0(11') + \int \mathrm{d}(22')\, G_0(12)\Sigma(22')G(2'1')$$



How to compute G?

The Dyson equation

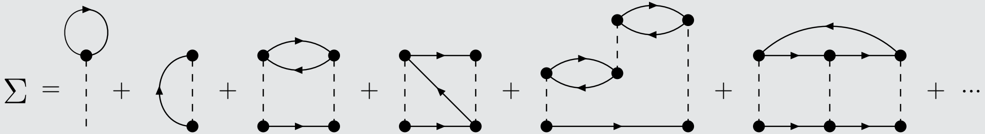
$$G(11') = G_0(11') + \int \mathrm{d}(22')\, G_0(12)\Sigma(22')G_0(2'1') + \int \mathrm{d}(22'33')\, G_0(12)\Sigma(22')G_0(2'3)\Sigma(33')G_0(3'1') + \ldots$$

How to compute G?

The Dyson equation

$$G(11') = G_0(11') + \int \mathrm{d}(22')\, G_0(12)\Sigma(22')G(2'1')$$

An exact expression for the self-energy



Another exact formalism

Self-consistent set of equations

$$G(11') = G_0(11') + G_0(12)\Sigma(22')G(2'1')$$

$$\Sigma(11') = \Sigma_{\mathrm{H}}(11') + \mathrm{i}\int \mathrm{d}(22'33')\ V(12;2'3)G(2'3')\Gamma(3'3;1'2)$$

$$\Gamma(12;1'2') = \delta(12')\delta(1'2) + \int \mathrm{d}(33'44')\,\frac{\delta\Sigma(11')}{\delta G(33')}G(34)G(4'3')\Gamma(42;4'2')$$

A few iterations

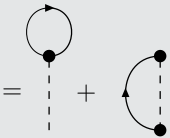
Initial condition

$$\Sigma^{(0)}(11') = 0 \qquad \Rightarrow \qquad \frac{\delta\Sigma^{(0)}(11')}{\delta G(33')} = 0$$

First iteration

$$\Gamma^{(1)}(12;1'2') = \delta(12')\delta(1'2') \qquad \Sigma^{(1)}(11') = \Sigma_{\mathrm{H}}(11') + \mathrm{i}\int \mathrm{d}(22')\ V(12;2'1')G(2'2)$$

$$\Sigma =$$



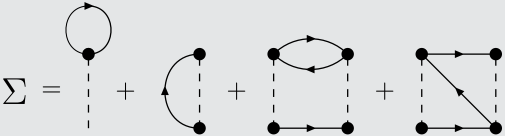
A few iterations

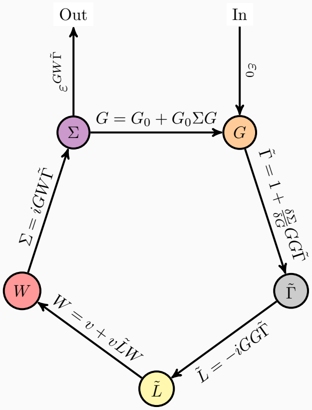
Second iteration

$$\frac{\delta\Sigma^{(1)}(11')}{\delta G(33')} = V(13';31') - V(13';1'3) = \bar{V}(13';31')$$

$$\Gamma^{(2)}(12;1'2') = \delta(12')\delta(1'2) + \int \mathrm{d}(33'44')\, \frac{\delta\Sigma^{(1)}(11')}{\delta G(33')} G(34)G(4'3')\Gamma^{(1)}(42;4'2')$$

$$\Sigma^{(2)}(11') = \Sigma_{\mathrm{Hx}}(11') + \mathrm{i}\int \mathrm{d}(22'33'44')\;\; V(12;2'3)G(2'3')\bar{V}(3'4';41')G(42)G(34)$$





Hedin, Phys Rev 139 (1965) A796

The wonderful equations of Hedin

$$G(11') = G_0(11') + \int G_0(12)\Sigma(22')G(2'1')\,\mathrm{d}(34)$$

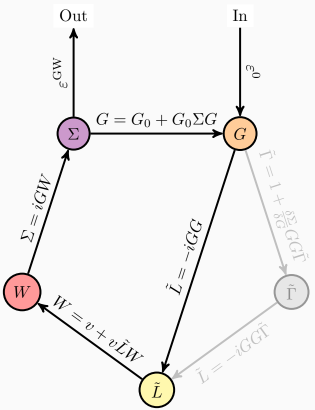
$$\underbrace{\tilde{\Gamma}(12;1'2')}_{\text{vertex}} = \delta(12')\delta(1'2) + \int \frac{\delta\Sigma_{\text{xc}}(11')}{\delta G(33')} G(34)G(4'3')\tilde{\Gamma}(42;4'2')$$

$$\underbrace{\tilde{L}(12;1'2')}_{\text{polarizability}} = -\mathrm{i} \int G(13)G(3'1')\tilde{\Gamma}(32;3'2')\,\mathrm{d}(33')$$

$$\underbrace{W(12;1'2')}_{\text{screening}} = V(12;1'2') + \int W(14;1'4')\tilde{L}(3'4';34)V(23;2'3')$$

$$\underbrace{\Sigma_{\text{xc}}(12)}_{\text{self-energy}} = \mathrm{i} \int G(33')W(12';32)\tilde{\Gamma}(3'2;1'2')\,\mathrm{d}(22'33')$$

Hedin's Square



Hedin, Phys Rev 139 (1965) A796

The *GW* approximation

$$G(11') = G_0(11') + \int G_0(12)\Sigma(22')G(2'1')\,\mathrm{d}(34)$$

$$\underbrace{\tilde{\Gamma}(12;1'2')}_{\text{vertex}} = \delta(12')\delta(1'2)$$

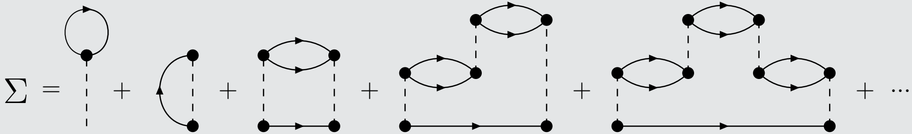
$$\underbrace{\tilde{L}(12;1'2')}_{\text{polarizability}} = -\mathrm{i}G(12')G(21')$$

$$\underbrace{W(12;1'2')}_{\text{screening}} = V(12;1'2') + \int W(14;1'4')\tilde{L}(3'4';34)V(23;2'3')$$

$$\underbrace{\Sigma_{\mathrm{xc}}(12)}_{\text{self-energy}} = \mathrm{i}\int G(32')W(12';31')\,\mathrm{d}(2'3)$$

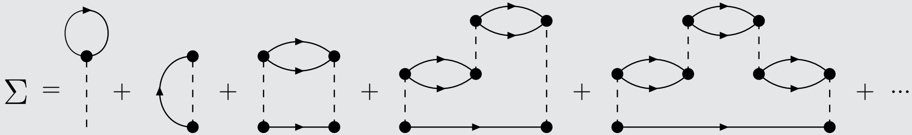
Diagrammatic content of the *GW* approximation

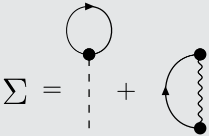
The *GW* resummation



Diagrammatic content of the *GW* approximation

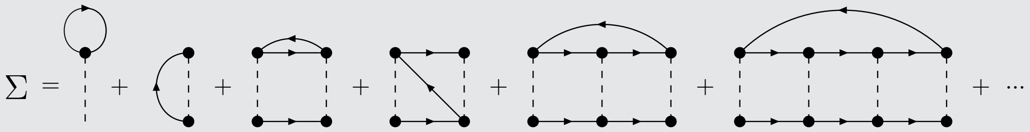
The *GW* resummation





Some other resummation-based self-energies

Particle-particle T-matrix



Some other resummation-based self-energies

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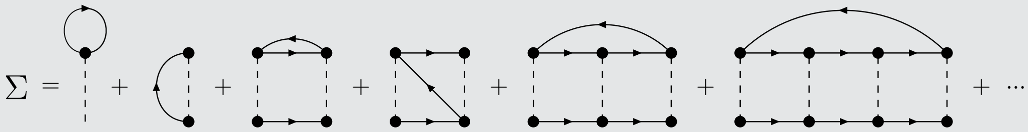
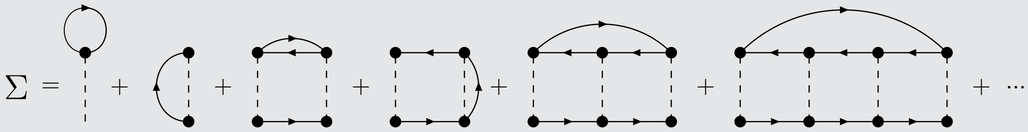


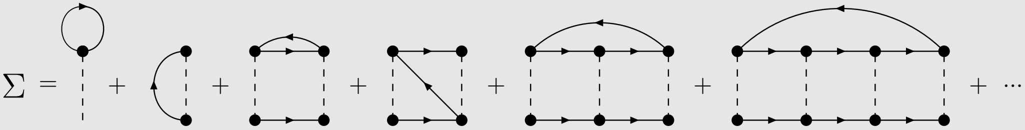
Electron-hole T-matrix



Romaniello, Bechstedt and Reining, Phys. Rev. B 85 (2012) 155131

Some other resummation-based self-energies

Particle-particle T-matrix

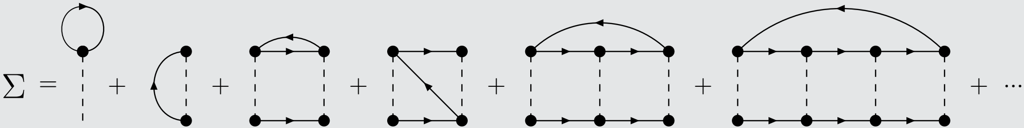


../diagrams/gtpp_square.pdf

What's missing with respect to the GW self-energy?

Some other resummation-based self-energies

Particle-particle T-matrix



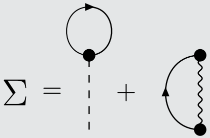
../diagrams/gtpp_square.pdf

What's missing with respect to the GW self-energy?

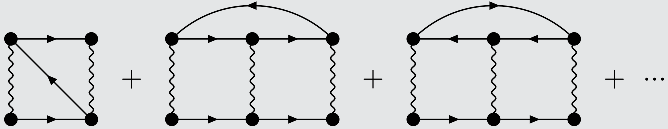
A systematic way to go beyond!

Vertex corrections to GW self-energy

First iteration



Second iteration



Mejuto-Zaera and Vlček, Phys. Rev. B 106 (2022) 165129

A new closed-set of equations for G

Key step of the derivation

Self-energy and equation of motion

$$\Sigma(11') = -\mathrm{i} \int \mathrm{d}(33'44')\, V(13;4'3')\, G_2(4'3';43)\, G^{-1}(41')$$

Two-body Green's function

Key step of the derivation

Self-energy and equation of motion

$$\Sigma(11') = -\mathrm{i} \int \mathrm{d}(33'44')\, V(13;4'3')\, G_2(4'3';43)\, G^{-1}(41')$$

Two-body Green's function

The Schwinger relation

$$G_2(12;1'2') = -\left.\frac{\delta G(11';[U])}{\delta\, U^{\text{eh}}(2'2)}\right|_{U=0} + G(11')G(22')$$

External potential

Key step of the derivation

Self-energy and equation of motion

$$\Sigma(11') = -\mathrm{i} \int \mathrm{d}(33'44')\, V(13;4'3')\, G_2(4'3';43)\, G^{-1}(41')$$

Two-body Green's function

The Schwinger relation

$$G_2(12;1'2') = -\left.\frac{\delta G(11';[U])}{\delta\, U^{\text{eh}}(2'2)}\right|_{U=0} + G(11')G(22')$$

External potential

The external potential

$$\hat{\mathcal{U}}(t_1) = \int \mathrm{d}(\mathbf{x}_1 \mathbf{x}_{1'} t_1')\, \hat{\psi}^\dagger(\mathbf{x}_1) U^{\text{eh}}(11') \hat{\psi}(\mathbf{x}_{1'})$$

Alternative Schwinger

Another external potential ...

$$\hat{\mathcal{U}}(t_1) = \frac{1}{2}\left(\int \mathrm{d}(\mathbf{x}_1\mathbf{x}_{1'}t_1')\,\hat{\psi}(\mathbf{x}_1)U^{\mathrm{hh}}(11')\hat{\psi}(\mathbf{x}_{1'}) + \int \mathrm{d}(\mathbf{x}_1\mathrm{d}\mathbf{x}_{1'}t_1')\,\hat{\psi}^\dagger(\mathbf{x}_{1'})U^{\mathrm{ee}}(11')\hat{\psi}^\dagger(\mathbf{x}_{1'})\right)$$

Alternative Schwinger

Another external potential ...

$$\hat{\mathcal{U}}(t_1) = \frac{1}{2}\left(\int \mathrm{d}(\mathbf{x}_1\mathbf{x}_{1'}t_1')\,\hat{\psi}(\mathbf{x}_1)U^{\mathrm{hh}}(11')\hat{\psi}(\mathbf{x}_{1'}) + \int \mathrm{d}(\mathbf{x}_1\mathrm{d}\mathbf{x}_{1'}t_1')\,\hat{\psi}^{\dagger}(\mathbf{x}_{1'})U^{\mathrm{ee}}(11')\hat{\psi}^{\dagger}(\mathbf{x}_{1'})\right)$$

...leading to an alternative Schwinger relation

Anomalous propagator

$$G_2(12;1'2') = -2\left.\frac{\delta\, G^{\mathrm{ee}}(1'2')}{\delta U^{\mathrm{hh}}(12)}\right|_{U=0}$$

Description of a non-number conserving Hamiltonian

Anomalous propagators

$$G^{\text{hh}}(11') = (-\mathrm{i}) \langle \Psi_0 | \hat{T} \left[\hat{\psi}(1) \hat{\psi}(1') \right] | \Psi_0 \rangle \qquad G^{\text{ee}}(11') = (-\mathrm{i}) \langle \Psi_0 | \hat{T} \left[\hat{\psi}^\dagger(1) \hat{\psi}^\dagger(1') \right] | \Psi_0 \rangle$$

Description of a non-number conserving Hamiltonian

Anomalous propagators

$$G^{\mathrm{hh}}(11') = (-\mathrm{i}) \langle \Psi_0 | \hat{T} \left[\hat{\psi}(1) \hat{\psi}(1') \right] | \Psi_0 \rangle \qquad G^{\mathrm{ee}}(11') = (-\mathrm{i}) \langle \Psi_0 | \hat{T} \left[\hat{\psi}^{\dagger}(1) \hat{\psi}^{\dagger}(1') \right] | \Psi_0 \rangle$$

Nambu formalism and the Gorkov propagator

$$\boldsymbol{G}(11') = (-i) \langle \Psi_0 | \hat{T} \left[\begin{pmatrix} \hat{\psi}(1)\hat{\psi}^{\dagger}(1') & \hat{\psi}(1)\hat{\psi}(1') \\ \hat{\psi}^{\dagger}(1)\hat{\psi}^{\dagger}(1') & \hat{\psi}^{\dagger}(1)\hat{\psi}(1') \end{pmatrix} \right] | \Psi_0 \rangle = \begin{pmatrix} G^{\mathrm{he}}(11') & G^{\mathrm{hh}}(11') \\ G^{\mathrm{ee}}(11') & G^{\mathrm{eh}}(11') \end{pmatrix}.$$

The particle-particle Hedin's equations

A new closed set of equations

$$G(12) = G_0(12) + \int G_0(13)\Sigma(34)G(42)\,\mathrm{d}(34)$$

$$\tilde{\Gamma}(12;1'2') = \frac{1}{2}\left(\delta(1'2)\delta(2'1) - \delta(1'1)\delta(2'2)\right) - \left.\frac{\delta\Sigma_{\mathrm{c}}^{\mathrm{ee}}(1'2')}{\delta G^{\mathrm{ee}}(33')}\right|_{U=0} G(43)G(4'3')\tilde{\Gamma}(12;44')$$

$$\tilde{K}(12;1'2') = \mathrm{i}G(31')G(3'2')\tilde{\Gamma}(12;33')$$

$$T(12;1'2') = -\bar{V}(12;1'2') - T(12;33')\tilde{K}(33';44')\bar{V}(44';1'2')$$

$$\Sigma_{\mathrm{Hxc}}(11') = \mathrm{i}G(2'2)T(12;33')\tilde{\Gamma}(33';2'1')$$

The particle-particle vertex

The T-matrix as a first approximation

$$G(12) = G_0(12) + \int G_0(13)\Sigma(34)G(42)\,\mathrm{d}(34)$$

$$\tilde{\Gamma}(12;1'2') = \frac{1}{2}\left(\delta(1'2)\delta(2'1) - \delta(1'1)\delta(2'2)\right)$$

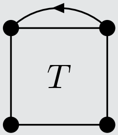
$$\tilde{K}(12;1'2') = \frac{\mathrm{i}}{2}\left(G(12')G(21') - G(22')G(11')\right)$$

$$T(12;1'2') = -\bar{V}(12;1'2') - T(12;33')\tilde{K}(33';44')\bar{V}(44';1'2')$$

$$\Sigma_{\text{Hxc}}(11') = \mathrm{i}G(2'2)T(12;1'2')$$

Vertex corrections to *GT* self-energy

First iteration

$\Sigma =$ 

Vertex corrections to GT self-energy

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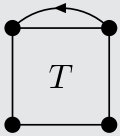
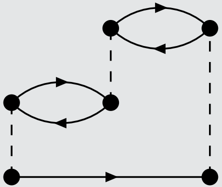
$\Sigma =$ 

Second iteration



Conclusion and perspectives

Conclusion and perspectives

Conclusions

- A new set of equations has been derived for the one-body propagator
- The pp T-matrix self-energy has no second-order term and the third order term might be really expensive
- Can we couple W and T thanks to the Nambu formalism?

Conclusion and perspectives

Conclusions

- A new set of equations has been derived for the one-body propagator
- The pp T-matrix self-energy has no second-order term and the third order term might be really expensive
- Can we couple W and T thanks to the Nambu formalism?

Anomalous propagators are also useful for two-body equations

- Simple expression for the kernel of the particle-particle channel!
- Accuracy of the particle-particle Bethe-Salpeter for double ionization?
- “Spin-flip-like” strategy for neutral excited states?

Questions?

A brief look at two-body equations

Two-body Bethe-Salpeter equations

The electron-hole Bethe-Salpeter equation

$$L(12;1'2') = L_0(12;1'2') + \int \mathrm{d}(3456)\, L_0(14;1'3)\Xi^{\text{eh}}(36;45)L(52;62').$$

Two-body Bethe-Salpeter equations

The electron-hole Bethe-Salpeter equation

$$L(12;1'2') = L_0(12;1'2') + \int \mathrm{d}(3456)\, L_0(14;1'3)\Xi^{\text{eh}}(36;45)L(52;62').$$

The particle-particle Bethe-Salpeter equation

$$K(12;1'2') = K_0(12;1'2') - \int \mathrm{d}(3456)\, K(12;56)\Xi^{\text{pp}}(56;34)K_0(34;1'2')$$

What's the difference?

The two kernels

$$\Xi^{\text{eh}}(12;34) = \left. \frac{\delta \Sigma^{\text{eh}}(13)}{\delta G^{\text{eh}}(42)} \right|_{U=0}$$

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The two kernels

$$\Xi^{\text{eh}}(12;34) = \left.\frac{\delta\Sigma^{\text{eh}}(13)}{\delta G^{\text{eh}}(42)}\right|_{U=0}$$

$$\int \mathrm{d}(3'44')\, G(24)\Xi^{\text{pp}}(34;3'4')K(3'4';1'2') = \int \mathrm{d}(3'44')\, G(41')\Xi^{\text{eh}}(34';43')L(3'2;4'2')$$

Csanak, Taylor and Yaris, Adv. Atom. Mol. Phys. 7 (1971) 287-361

What's the difference?

The two kernels

$$\Xi^{\text{eh}}(12;34) = \left.\frac{\delta\Sigma^{\text{eh}}(13)}{\delta G^{\text{eh}}(42)}\right|_{U=0}$$

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Csanak, Taylor and Yaris, Adv. Atom. Mol. Phys. 7 (1971) 287-361

A new expression for the particle-particle kernel

$$\Xi^{\text{pp}}(12;34) = \left.\frac{\delta\Sigma^{\text{ee}}(34)}{\delta G^{\text{ee}}(12)}\right|_{U=0}$$